Cooling across a superfluid-normal interface in spinimbalanced fermions - An analogue of a dilution fridge ?

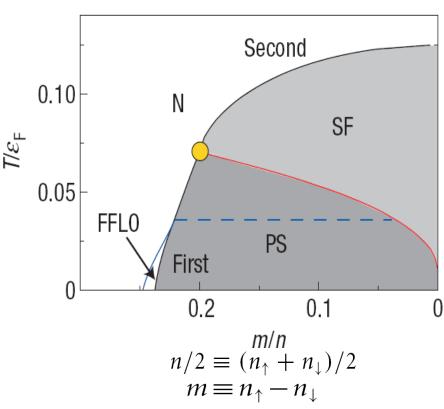
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Spin-imbalanced fermions close to unitarity

• The phase diagram of spin-imbalanced fermionic atoms interacting via a feshbach resonance shares interesting features with He₃-He₄ mixture.

 At higher temperatures there is a line of second order transitions between a paired superfluid and normal leading to a tri-critical point.

 Below the tri-critical point the transition turns first order and the system phase separates between partially polarized normal and an unpolarized superfluid.



Mean-field phase diagram at $1/k_Fa = -1$ Parish et. al., Nature Physics 3, 124 - 128 (2007)

Only qualitatively correct (except perhaps FFLO) but quantitatively inaccurate.

Experiments with spin-imbalanced fermions (MIT)

Zwierlein et. al., Nature (London) 442, 54 (2006)

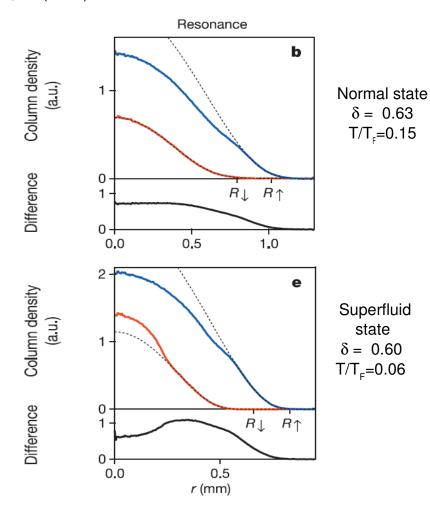
The polarized superfluid is unstable at high polarization.
What is critical polarization at equilibrium?

Polarization δ is defined to be

$$\delta = (N_{\uparrow} - N_{\downarrow})/(N_{\uparrow} + N_{\downarrow}))$$

• The critical polarization observed in the MIT group δ_c ≈ 0.7 (consistent with QMC predictions assuming Local Density approximation(LDA))

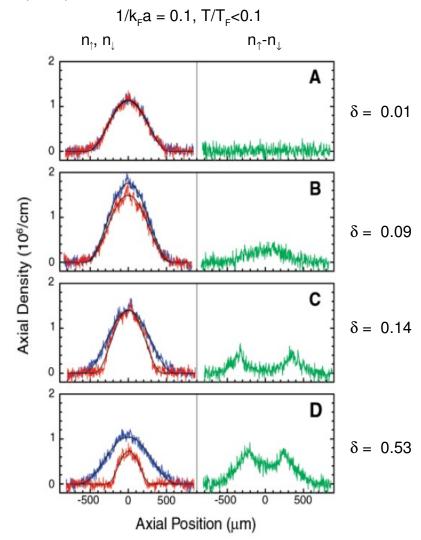
Trap aspect ratio ≈ 5



Experiments with spin-imbalanced fermions (Rice)

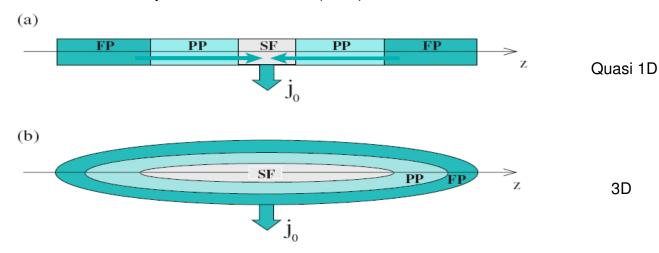
Partridge et. al., Science 311, 503 (2006)

- The critical polarization observed in the Rice group $\delta_{\rm c}\approx 0.9$ (apparent violation of Clogston limit)
- Trap aspect ratio ≈ 30
- The spin densities violate LDA
- Both experiments report temperatures which are low enough for δ_c to be close to the T=0 value.



Non-equilibrium scenario in quasi 1D

Meera Parish and David Huse Phys. Rev. A 80, 063605 (2009)



- For a highly elongated trap, evaporation occurs from the SF phase predominantly removing
 atoms at low T.
- This generates a chemical potential gradient producing a spin current along the axial direction.
- But the spin-gap of the SF blocks spin current across the SF-N interface thus effectively depolarizing the SF.

Metastable phase

Y. A. Liao, et. al., arXiv:1105.6369

Axial density (10⁵ cm⁻¹)

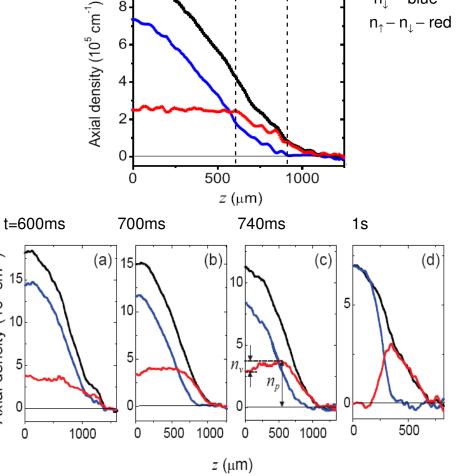
 $t=3.4 \text{ s}; T/T_{E}=0.09$

Ramping the trap depth exponentially from U=160μK to $2.2\mu K$ over 3.4s.

Consistent with predictions based on LDA.

Ramping the trap depth exponentially from U=160μK to 0.74μK over 1s (more aggressive evaporation).

A dip in the central density difference violating LDA.



n_↑ – black

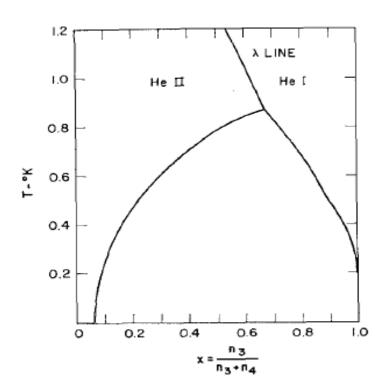
 n_{\perp} – blue

 $n_{\uparrow}\!-n_{\downarrow}\!-red$

He₃ – He₄ mixture phase diagram

- At high temperatures, there is the λ-transition between superfluid and normal phase of He₄.
 Critical temperature is suppressed by the presence of fermionic He₃.
- At low temperatures, the system phase separates into a He₃ rich phase ("liquid") and a dilute solution of He₃ in He₄ superfluid ("vapour").

Even at T=0 there is a finite concentration of He₃ in He₄
due to the degeneracy of the He₃ fermions and interaction effects.

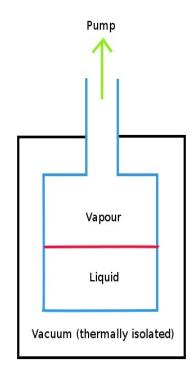


Basic working-principle of a dilution fridge

For the He₃-He₄ mixture, the "*vapour*" phase corresponds to the dilute solution of He₃ in superfluid He₄ which is effectively at T=0.

• The "*liquid*" phase corresponds to the concentrated He₃ with very little He₄.

Evaporating the He₃ atoms from the "*liquid*" phase results in cooling the "*liquid*" and the *interface* separating the 2 phases.



Transport across SF-N interface

 Under conditions when the system is driven (via evaporative cooling) a steady state can be established where

$$\mu_{\uparrow}^N + \mu_{\downarrow}^N \neq 2\mu_{SF}$$

due the blockage of spin transport across the interface.

 Another consequence of such a mechanism, would inhibit the transmission of heat thus resulting in a gradient in temperature as well.

$$T_N \neq T_{SF}$$

